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Outline

Harmonic
Maps:
Riemannian
case

Generalization:
Pseuro-
Riemannian
case

Summary

Harmonic maps between pseudo-Riemannian surfaces

A. Fotiadis (coauthor C. Daskaloyannis)

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Outline

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1 Harmonic Maps: Riemannian case

2 Generalization: Pseuro-Riemannian case

3 Summary

Harmonic maps

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Summary

M, N Riemannian manifolds, $u: M \rightarrow N$

$$e(u) = \text{Tr}_g(u^*(h)) = g^{ij}(x)h_{\alpha\beta}(u(x))\partial_i u^\alpha \partial_j u^\beta.$$

$$E(u) = \int_M e(u) dv_M,$$

Euler-Lagrange equations of the energy functional are a nonlinear system of elliptic PDEs of second order.

Solutions: Harmonic Maps

Riemann surfaces

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Mapping between surfaces: Isothermal coordinates...

$$g = e^{f(z, \bar{z})} |dz|^2, h = e^{F(u, \bar{u})} |du|^2$$

$$E(u) = \int e^{F(u, \bar{u})} (|\partial_z u_z|^2 + |\partial_{\bar{z}} u|^2) dx dy,$$

Harmonic map equations:

$$\partial_{z\bar{z}} u + \partial_u F(u, \bar{u}) \partial_z u \partial_{\bar{z}} u = 0.$$

Integrability

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$u = R + iS$ harmonic iff

$$\partial_{\bar{z}} \left(e^{F(u, \bar{u})} \partial_z u \partial_z \bar{u} \right) = 0$$

Consider locally a conformal change of coordinates s.t.

$$e^{F(u, \bar{u})} \partial_z u \partial_z \bar{u} = 1$$

Change of coordinates

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If

$$e^{F(u, \bar{u})} \partial_z u \partial_{\bar{z}} \bar{u} = e^{-\lambda(z)}$$

then

$$\zeta = \xi + i\eta = \int e^{-\lambda(z)/2} dz$$

Parametrization

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Set

$$e^v = e^{\frac{F}{2}} \partial_z u \text{ and } v = \omega + i\theta.$$

Then

$$e^v = e^{\frac{F}{2}} \partial_z u$$

$$e^{-\bar{v}} = e^{\frac{F}{2}} \partial_{\bar{z}} u$$

Observation

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solutions of Beltrami equation

$$\frac{\partial_{\bar{z}}u}{\partial_z u} = e^{-2\omega}$$

Compatibility conditions

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Since $\partial_{z\bar{z}}^2 v = \partial_{\bar{z}z}^2 v$ we get:

$$\partial_x \omega - \partial_y \theta = e^{-\frac{F(R,S)}{2}} \sinh \omega (\partial_R F \cos \theta + \partial_S F \sin \theta)$$

$$\partial_x \theta + \partial_y \omega = e^{-\frac{F(R,S)}{2}} \cosh \omega (\partial_S F \cos \theta - \partial_R F \sin \theta)$$

sinh-Gordon

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Since $\partial_{xy}^2 \theta = \partial_{yx}^2 \theta$ we get:

$$\Delta \omega = 2e^{-F} \Delta F \sinh \omega \cosh \omega = -2K_N \sinh 2\omega$$

Bäcklund transformation

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$$\begin{aligned}\Delta\omega &= -2K_N \sinh(2\omega), \text{ Bäcklund transformation} \\ i\Delta\theta &= 2e^{-F} \left((F_u^2 - F_{uu}) e^{2i\theta} - (F_{\bar{z}}^2 - F_{\bar{z}\bar{z}}) e^{-2i\theta} \right)\end{aligned}$$

Bäcklund transformation

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Let

$$e^F = \frac{1}{S^2}.$$

$$\partial_x \omega - \partial_y \theta = -2 \sinh \omega \sin \theta$$

$$\partial_x \theta + \partial_y \omega = -2 \cosh \omega \cos \theta$$

Bäcklund transformation:

$$\Delta \omega = -2 \sinh (2\omega) \leftrightarrow \Delta \theta = -2 \sin (2\theta)$$

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Constant Curvature: Uniform Approach

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Summary

We have to solve sinh-Gordon

$$\Delta\omega = -2K_N \sinh 2\omega$$

and Beltrami

$$\frac{\partial_{\bar{z}}u}{\partial_z u} = e^{-2\omega}$$

Assume

$$\omega = \omega(ax + by)$$

one soliton solution...

Elliptic functions

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$$x = \int_0^\phi \frac{1}{\sqrt{1 - m \sin^2 t}} dt$$

$$\operatorname{sn}(x, m) = \sin \phi$$

$$\operatorname{cn}(x, m) = \cos \phi$$

$$\operatorname{dn}(x, m) = \sqrt{1 - m \sin^2 x}$$

$$\Pi(n, \phi, m) = \int_0^\phi \frac{1}{(1 - n \sin^2 t) \sqrt{1 - m \sin^2 t}} dt$$

1-soliton solutions

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Summary

Solve sinh-Gordon equation when $K_N = \text{const}$ and
 $\omega = \omega(ax + by)$:

$$\Delta\omega = -2K_N \sinh 2\omega$$

via elliptic integrals

1-soliton solutions

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$$R - R_0 = X - X_0 + i \frac{\tan \tau}{c} \left(\Pi \left(\frac{1}{\cos^2 \tau}, iw, -\frac{4K_N}{c\rho^2} \right) - \Pi_0 \right)$$

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$$S - S_0 = \frac{1}{\sqrt{c} \sqrt{1 + \frac{4K_N}{c\rho^2} \cos^2 \tau}} \left(\operatorname{arctanh} \frac{w'(Y)}{\sqrt{c} \sqrt{1 + \frac{4K_N}{c\rho^2} \cos^2 \tau}} - a_0 \right)$$

$$e^{F(S)} = \frac{\cos^2 \tau}{K_N} \left(1 + \frac{4K_N}{c\rho^2} \right) \frac{1}{\cosh \left(\sqrt{c \left(1 + \frac{4K_N}{c\rho^2} \cos^2 \tau \right)} (S - S_0) + \Sigma_0 \right)}$$

where

$$\tanh \omega = \frac{1}{\sqrt{m}} \operatorname{sn} \left(\sqrt{Cm} (Y - Y_0) + v_0 \middle| \frac{1}{m} \right) \dots$$

Recover example in JDG

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$$R = x + a \int_0^y \sin^2 S(t) dt,$$

$$S = \cot^{-1} z(y), y \in [0, \frac{\pi}{2}]$$

$$\int_0^{z(y)} \frac{dz}{\sqrt{z^4 + c^2 z^2 + b^2}} = \frac{\pi}{2} - y, c^2 = 1 + b^2 + a^4, y \in [0, \frac{\pi}{2}]$$

Recover example in JDG

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$$\int_0^{\infty} \frac{dz}{\sqrt{z^4 + c^2 z^2 + b^2}} = \frac{\pi}{2}$$

$$a^2 \int_0^{\infty} \frac{dz}{(1+z^2)\sqrt{z^4 + c^2 z^2 + b^2}} = \frac{\beta}{2}$$

and extend the functions R, S by symmetry in the $y = \frac{\pi}{2}$ axis...
Recover examples: Shi-Tam-Wan, Wolf, Li-Tam

Pseudo-Riemann surfaces

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Mapping between surfaces: Isothermal coordinates...

$$g = e^{f(v,w)} dv dw = e^{f(x,y)} (dx^2 - \epsilon^2 dy^2),$$

where $v = x + \epsilon y$, $w = x - \epsilon y$, $\epsilon = 1, i$

$$h = e^{F(V,W)} dV dW = e^{F(R,S)} (dR^2 - \delta^2 dS^2),$$

where

$v = x + \epsilon y$, $w = x - \epsilon y$, $V = R + \delta S$, $W = R - \delta S$, $\delta = 1, i$

Harmonic maps

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M, N pseudo-Riemann surfaces, $u: M \rightarrow N$

In isothermal coordinates:

$$E(u) = \frac{1}{2} \int (R_x^2 - \frac{1}{\epsilon^2} R_y^2 - \delta^2 S_x^2 + \frac{\delta^2}{\epsilon^2} S_y^2) dx dy$$

Euler-Lagrange equations of the energy functional are a nonlinear system of elliptic or hyperbolic PDEs of second order.
Solutions: Harmonic (or Wave) Maps

Notation

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$$V_v = \frac{1}{2}(V_x + \frac{1}{\epsilon}V_y) = \frac{1}{2}(R_x + \frac{1}{\epsilon}R_y + \delta S_x + \frac{\delta}{\epsilon}S_y),$$

$$V_w = \frac{1}{2}(V_x - \frac{1}{\epsilon}V_y) = \frac{1}{2}(R_x - \frac{1}{\epsilon}R_y + \delta S_x - \frac{\delta}{\epsilon}S_y),$$

$$W_v = \frac{1}{2}(W_x + \frac{1}{\epsilon}W_y) = \frac{1}{2}(R_x + \frac{1}{\epsilon}R_y - \delta S_x - \frac{\delta}{\epsilon}S_y),$$

$$W_w = \frac{1}{2}(W_x - \frac{1}{\epsilon}W_y) = \frac{1}{2}(R_x - \frac{1}{\epsilon}R_y - \delta S_x + \frac{\delta}{\epsilon}S_y).$$

Harmonic maps

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A map $u = (R, S)$ is harmonic if it satisfies the system

$$V_{vw} + F_V(V, W) V_v W_w = 0$$

$$W_{vw} + F_W(V, W) W_v W_w = 0,$$

where $V = R + \delta S$, $W = R - \delta S$.

Integrability

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Summary

$u = R + iS$ harmonic iff

$$\left(e^{F(V,W)} V_v W_v \right)_w = 0$$

and

$$\left(e^{F(V,W)} V_w W_w \right)_v = 0$$

Consider locally a conformal change of coordinates s.t.

$$\left(e^{F(V,W)} V_v W_v \right) = 1, \left(e^{F(V,W)} V_w W_w \right) = 1$$

Change of coordinates

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If

$$\left(e^{F(V,W)} V_v W_v \right) = \Lambda(v), \quad \left(e^{F(V,W)} V_w W_w \right) = M(w)$$

then

$$\zeta = \int \frac{1}{\sqrt{\Lambda(v)}} dv, \quad \eta = \int \frac{1}{\sqrt{M(w)}} dw$$

Parametrization

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Harmonic maps equations are equivalent to...

$$R_x R_y - \delta^2 S_x S_y = 0$$

and

$$\left(R_x^2 + \frac{1}{\epsilon^2} R_y^2 \right) - \delta^2 \left(S_x^2 + \frac{1}{\epsilon^2} S_y^2 \right) = 4e^{-F(R,S)}.$$

Parametrization

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$$R_x = 2e^{-\frac{F}{2}} \cosh \Omega \cosh \Theta$$

$$R_y = 2\epsilon e^{-\frac{F}{2}} \sinh \Omega \sinh \Theta$$

$$S_x = \frac{2}{\delta} e^{-\frac{F}{2}} \cosh \Omega \sinh \Theta$$

$$S_y = \frac{2\epsilon}{\delta} e^{-\frac{F}{2}} \sinh \Omega \cosh \Theta,$$

where

$$\Omega = \frac{\delta}{\epsilon} \omega, \quad \Theta = \delta \theta$$

Compatability conditions

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Summary

Since $V_{vw} = V_{wv}$ and $W_{vw} = W_{wv}$ we get:

$$\begin{aligned}\Omega_w + \Theta_w &= \frac{F_w}{2} - F_V V_w, \\ -\Omega_v + \Theta_v &= \frac{F_v}{2} - F_V V_v.\end{aligned}$$

sine-Gordon type equation

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Since $\partial_{xy}^2 \Theta = \partial_{yx}^2 \Theta$ we get:

$$\Omega_{xx} - \epsilon^2 \Omega_{yy} = -2K_N \sinh 2\Omega.$$

where

$$K_N = -2F_{VW} e^{-F} = -\frac{1}{2} (F_{RR} - \delta^2 F_{SS}) e^{-F(R,S)},$$

and $\Omega = \frac{\delta}{\epsilon} \omega$

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Observation

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solutions of Beltrami equations

$$\frac{V_w}{V_v} = e^{-2\Omega},$$

$$\frac{W_w}{W_v} = e^{2\Omega},$$

Bäcklund transformation

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Summary

$$\begin{aligned}\Omega_{vw} &= -\frac{K_N}{2} \sinh(2\Omega), \text{ Bäcklund transformation} \\ 2\Theta_{vw} &= e^{-F} \left((F_V^2 - F_{VV}) e^{2\Theta} - (F_W^2 - F_{WW}) e^{-2\Theta} \right)\end{aligned}$$

Bäcklund transformation

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Summary

Let

$$e^F = \frac{1}{(aR + \delta^2 bS)^2}.$$

If $\delta = \epsilon = 1$, $a = 1$, $b = 0$ then auto-Bäcklund of

$$\omega_{xx} - \omega_{yy} = 2 \sinh(2\omega),$$

If $\delta = i$, $\epsilon = 1$, $a = 1$, $b = 0$ then auto-Bäcklund of

$$\omega_{xx} - \omega_{yy} = 2 \sin(2\omega).$$

Constant Curvature: Uniform Approach

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We have to solve

$$\Omega_{xx} + \frac{1}{\epsilon^2} \Omega_{yy} = -\frac{K_N}{2} \sinh 2\Omega, \text{ where } \Omega = \frac{\delta}{\epsilon} \omega.$$

and Beltrami equations

$$\frac{V_w}{V_v} = e^{-2\Omega},$$

$$\frac{W_w}{W_v} = e^{2\Omega},$$

Assume

$$\omega = \omega(ax + by)$$

one soliton solution...

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Interesting Problems

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Summary

1. Find new solutions to the sinh-Gordon equation and the corresponding harmonic maps.
2. The generalized sine-Gordon equation is interesting on its own right. Construct new families of solutions.
3. Multidimensional cases (under current investigation)

Publications

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- Beltrami equation for the harmonic diffeomorphisms between surfaces. A. Fotiadis and C Daskaloyannis, Nonlinear Analysis (2022).
- Unified formalism for harmonic maps between Riemann or Lorentz surfaces. A. Fotiadis and C Daskaloyannis, Accepted in Journal of Geometry and Physics
- New examples of harmonic maps to the hyperbolic plane via Backlund transformation. G. Polychrou, E. Papageorgiou, A. Fotiadis, C. Daskaloyannis. Revista Matematica Complutense (2023)

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Summary

- Harmonic maps between riemannian surfaces
- Generalization
- Open problems

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Thank you very much for your attention!