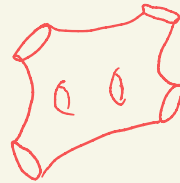



Moduli spaces of flat connections and the Atiyah-Bott classes

Yiannis Loizides

George Mason University

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Moduli spaces of flat connections on surfaces

G - compact Lie group with $\pi_0 G = \pi_1 G = \{1\}$ (e.g. SU_N)

\mathfrak{g} - Lie(G), with positive definite integral $k \in \text{Sym}^2(\mathfrak{g}^*)^{\mathfrak{g}}$ "level"

V - faithful representation of G (e.g. \mathbb{C}^N)

Σ - compact connected oriented surface with $\partial\Sigma \cong S^1$ parametrized

A connection (on $\Sigma \times V$) is an operator

$$d_A = d + A : \Omega^\bullet(\Sigma, V) \rightarrow \Omega^{\bullet+1}(\Sigma, V)$$

where $A \in \Omega^1(\Sigma, \mathfrak{g})$

$$A(\Sigma) = \left\{ d_A \mid A \in \mathcal{A}'(\Sigma, \mathfrak{g}) \right\}$$

curvature $F_A = d_A^2 = dA + \frac{1}{2} [A, A] \in \mathcal{A}^2(\Sigma, \mathfrak{g})$

$\text{Map}(\Sigma, G) \overset{C}{\cong} A(\Sigma)$ by gauge transformations

$$g: d_A \longmapsto g \circ d_A \circ g^{-1} = d_{A^g}, \quad A^g = gAg^{-1} - dg g^{-1}$$

$$F_{A^g} = g F_A g^{-1}$$

d_A is flat if its curvature $F_A = 0$

$$\mathcal{A}_{\text{flat}}(\Sigma) = \{ d_A \mid A \in \Omega^1(\Sigma, \mathfrak{g}), d_A^2 = 0 \}$$

Action of $\text{Map}_\partial(\Sigma, G) = \{ g \in \text{Map}(\Sigma, G) \mid g|_{\partial\Sigma} = 1 \}$ free and proper.

The moduli space of flat connections with framing along $\partial\Sigma$ is

$$\mathcal{M}_\Sigma = \mathcal{A}_{\text{flat}}(\Sigma) / \text{Map}_\partial(\Sigma, G)$$

The moduli space of flat connections with framing along $\partial\Sigma$ is

$$\mathcal{M}_\Sigma = \mathcal{A}_{\text{flat}}(\Sigma) / \text{Map}_\partial(\Sigma, G) \quad (*)$$

- smooth infinite-dim. manifold
- weakly symplectic $(*)$ is a symplectic quotient, $\omega_{AB}(a, b) = \int_\Sigma \ell(a, b)$, [Atiyah-Bott], [Donaldson]
- Hamiltonian action of $\text{Map}(\Sigma, G) / \text{Map}_\partial(\Sigma, G) \cong LG = \text{Map}(S^1, G)$

proper moment map $\mu: \mathcal{M}_\Sigma \longrightarrow \Omega^1(S^1, \mathfrak{g}) \subseteq \text{Log}^*$

$$[d_A] \longmapsto A|_{\partial\Sigma}$$

Hamiltonian loop group spaces

A proper Hamiltonian LG-space is an infinite-dim. weakly symplectic manifold (M, ω) carrying a Hamiltonian action of $LG = \text{Map}(S^1, G)$ with proper moment map $\mu: M \rightarrow \Omega^1(S^1, \mathfrak{g}) \subseteq \text{Log}^*$ [Meinrenken-Woodward]

$$\iota(X)\omega = -d \int_{S^1} \ell(\mu, X) \quad \forall X \in \text{Log} = \Omega^0(S^1, \mathfrak{g})$$

Example orbits of $LG \curvearrowright \Omega^1(S^1, \mathfrak{g})$ (coadjoint orbits)

Rmk [M-W] use Sobolev spaces of loops, and M is a Banach manifold.

Suppress this for simplicity.

Motivation

$$M_{\Sigma} = \mathcal{A}_{\text{flat}}(\Sigma) / \text{Map}_{\partial}(\Sigma, G) \xrightarrow{\mu} \Omega'(S^1, \mathfrak{g})$$

→ symplectic quotient $M_{\Sigma,0} = \mu^{-1}(0)/G$ is the moduli space of flat connections on $\Sigma \cup_{\partial\Sigma} \mathbb{D}$, known to algebraic geometers as the moduli space of semistable $G_{\mathbb{C}}$ -bundles

→ topology of $M_{\Sigma,0}$ can be studied via G -equivariant topology of M_{Σ}

Example Verlinde formula for the index of line bundles on $M_{\Sigma,0}$
(Jeffrey-Kirwan, Bismut-Labourie, Alekseev-Meinrenken-Woodward, ...)

→ in algebraic geometry, Teleman-Woodward proved a generalization of the Verlinde formula covering a broad collection of K -theory classes

K-theory

H - Hilbert space, $\dim(H) = \infty$

$$\mathcal{F} = \{ T \in \mathcal{B}(H) \mid T \text{ is Fredholm} \} \quad (\text{norm topology})$$

$$K(X) = [X, \mathcal{F}] \quad (\text{topologist's/non-compact support})$$

Rmk If X is loc. compact, $\mathcal{R}KK(X; C_0(X), C_0(X))$ in Kasparov's notation.

Rmk Variants with $T: H_1 \rightarrow H_2$.

Rmk G compact Lie group, $K_G(X)$.

$\bar{\partial}$ operators

Fix a complex structure on Σ .

$$d_A \rightsquigarrow (d_A)^{0,1} = \bar{\partial}_A : \Omega^0(\Sigma, V) \longrightarrow \Omega^{0,1}(\Sigma, V)$$

family $\{\bar{\partial}_A \mid d_A \in \mathcal{A}_{\text{flat}}(\Sigma)\}$ is gauge-equivariant, but does not define a K -class because

$$\partial\Sigma \neq \emptyset \implies \bar{\partial}_A \text{ not Fredholm}$$

Boundary conditions

On the complex disk \mathbb{D} , $\ker(\bar{\partial})$ is infinite dimensional

$$1, z, z^2, z^3, \dots$$

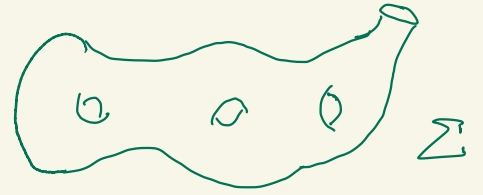
To cut down to finite dimensions, impose a boundary condition

$$f|_{\partial\mathbb{D}} \in B_{<0}(\frac{1}{i}\partial_{\theta}) = \left\{ \sum_{n<0} a_n e^{in\theta} \right\}$$

then $\ker(\bar{\partial}^b: \Omega^0(\mathbb{D}; B_{<0}(\frac{1}{i}\partial_{\theta})) \rightarrow \Omega^{0,1}(\mathbb{D})) = \emptyset$

On a general Σ , we do

the same: the operator



$$\bar{\mathcal{D}}_A^b : \Omega^0(\Sigma, V; B_{<0}(\frac{1}{i} \partial_{\theta})) \longrightarrow \Omega^{0,1}(\Sigma, V) \quad (*)$$

is Fredholm.

Rmk: This is essentially the Atiyah-Patodi-Singer boundary condition.

Note however that we take $B_{<0}(\frac{1}{i} \partial_{\theta})$ to be independent of A .

(*) Slightly cleaner if we use a spin str on Σ . We assume this below.

The family $\{ \bar{\partial}_A^b \mid d_A \in \mathcal{A}(\Sigma) \}$ is not gauge - equivariant:

$$B_{\leq 0} \left(\frac{1}{i} \partial_{\theta} \right) \xrightarrow{g \in \text{Map}(\Sigma, G)} B_{\leq 0} \left(g_{\partial} \circ \frac{1}{i} \partial_{\theta} \circ g_{\partial}^{-1} \right), \quad g_{\partial} = g|_{\partial \Sigma}$$

Preserved by $g: \Sigma \rightarrow G$ such that $g|_{\partial \Sigma}$ is constant, and $\frac{\text{Map}_{\text{const}}(\Sigma, G)}{\text{Map}_{\partial}(\Sigma, G)} \cong G$.

Thm [L] The family $\{ \bar{\partial}_A^b \}$ descends to a continuous family of Fredholm operators parametrized by $M_{\Sigma} = \frac{\mathcal{A}_{\text{flat}}(\Sigma)}{\text{Map}_{\partial}(\Sigma, G)}$.

Notation: $E^{\Sigma} V \in K_G(M_{\Sigma})$

The index homomorphism

M - proper Hamiltonian LG -space,
such as M_Σ

Next, define a homomorphism

$$\text{index}_l : K_T(M)_{gc} \longrightarrow R^{-\infty}(T) = \mathbb{Z}^\Lambda$$

T - max. torus of G
 $\Lambda = \text{Hom}(T, U(1))$

gc is a "growth condition" - to be explained

Thm [L] For $l \gg 0$, the coefficient of e^λ in $\text{index}_l(\Psi)$ computes a K -theoretic intersection pairing on a finite dimensional symplectic quotient of M .

A triple (X, ξ, f) where

X - finite even dim. T -equiv. Spin^c

$\xi \in K_{T,c}(X)$ (compactly supported K -theory)

$f: X \longrightarrow M$ T -equiv.

pairs with $K_T(M)$: $\langle [(X, \xi, f)], \psi \rangle = \text{index}_T^X(\xi \otimes f^* \psi) \in R(T)$

Definition of index_ℓ involves a non-compactly supported ξ .

\rightarrow pairing only defined for subring $K_T(M)_{gc}$

\rightarrow output in $R^{-\infty}(T)$

Construction of (X, ξ, f)

Thm [L-Meinrenken-Song] \exists $\dim(\mathfrak{g})$ -dim. submanifold $\mathfrak{t} \subseteq U \subseteq \Omega^1(S^1, \mathfrak{t})$
which is transverse to gauge orbits (canonical up to homotopy).

$$X = \mu^{-1}(U), \quad f = \text{inclusion}, \quad \xi = \mu^* \text{Thom}(U, \mathfrak{t}^*)$$

$$X \hookrightarrow \text{Norm}_{\text{LG}}(\mathfrak{t}^*) \stackrel{\substack{\cong \\ \text{finite} \\ \text{index}}}{=} N = \left\{ \begin{array}{c} \text{closed} \\ \text{geodesics} \\ \text{in } T \end{array} \right\} \cong T \times \Pi, \quad \Pi = \text{Hom}(S^1, T)$$

Thm [L-Meinrenken-Song] \exists canonical Spin^c structure on X induced
by a compatible almost complex structure on the symplectic manifold M .

Growth condition

$$N = \left\{ \begin{array}{l} \text{closed} \\ \text{geodesics} \\ \text{in } T \end{array} \right\}$$

Assume the $K_T(M)$ - class represented by a smooth T -equiv. family of operators with compact resolvent:

$$\begin{array}{ccc} \mathcal{E}^+ \supseteq \text{dom}(Q) & \xrightarrow{Q} & \mathcal{E}^- \\ & \searrow & \swarrow \\ & M & \end{array}$$

(\mathcal{E}, Q) satisfies the gc if

① $\mathcal{E}|_X$ is N -equiv., using the **norm topology** on $\text{Aut}(\mathcal{E}|_X)$

② \exists N -equiv. connection $\nabla^{\mathcal{E}|_X}$ preserving $\text{dom}(Q)|_X$ such that

$\nabla^{\mathcal{E}|_X} Q|_X$ is a **bounded** section of $T^*X \otimes \mathcal{B}(\mathcal{E}|_X)$

Rmk $\text{dom}(Q)|_X \in \mathcal{E}^+|_X$ not required to be N -invariant.

Finite rank N -equiv. vector bundles satisfy the g.c.

Thm [L] The family of Fredholm operators on \mathcal{M}_Σ induced by $\{ \bar{\partial}_A^b \mid d_A \in \mathcal{A}_{\text{flat}}(\Sigma) \}$ satisfies the g.c.

There are K -classes that do not satisfy the g.c.

Rmk The 'size' of $K_T(\mathcal{M})_{\text{g.c.}}$ in $K_T(\mathcal{M})$ remains unclear (to me).

Thm [L] There is a homomorphism

$$\text{index}_\ell: K_T(M)_{gc} \longrightarrow R^{-\infty}(T) = \mathbb{Z}^\wedge$$

that takes $[(\mathcal{E}, \mathcal{Q})]$ to the L^2 -index of $\mathbb{D}_{X, \xi}^d \hat{\otimes}_{\nabla^{\mathcal{E}}} 1 + 1 \hat{\otimes} \mathcal{Q}_X$

\mathcal{L} = prequantum line bundle on M , at level ℓ

$$\text{i.e. } c_1(\mathcal{L}) = [\omega] \in H^2(M), \quad \mathcal{L} \in \widehat{LG}^{(\ell)}$$

Rmk \mathcal{L} does not satisfy the gc.

Rmk Generalizes [L-Song], which treated finite rank N -equiv. vector bundles.

Example

$$\Sigma = \mathbb{D}, \quad \mathcal{M}_{\mathbb{D}} \cong LG/G \quad (\text{"smooth affine Grassmannian"})$$

$$X \sim \mathcal{M}_{\mathbb{D}}^{T \times S^1_{\text{rot}}} \cong \Pi = \text{Hom}(S^1, T) \xleftarrow{\ell} \Lambda = \text{Hom}(\Pi, \mathbb{Z})$$

↑ Π -orbit of trivial connection on \mathbb{D}

$$\begin{aligned} \text{index}_k(E^{\mathbb{D}}V) &= J \sum_{\eta \in \Pi} e^{\eta} \cdot \text{index}(\bar{\partial}_{\tilde{\eta}}^b) \quad , \quad J = \begin{array}{l} \text{Weyl} \\ \text{denominator} \\ \text{for } G \end{array} \\ &= J \sum_{\eta \in \Pi} e^{\eta} \langle d\chi_V, \eta \rangle \end{aligned}$$

where $\chi_V(u) = \text{Tr}_V(\rho_V(u))$, $u \in T$ character of (V, ρ_V)

$$\text{index}_k(E^{\mathbb{D}}V) = J \sum_{\eta \in \Pi} e^{\eta} \langle d\chi_V, \eta \rangle$$

Consider the generating series

$$e^{tE^{\mathbb{D}}V} = \sum_{m \geq 0} \frac{t^m}{m!} (E^{\mathbb{D}}V)^{\otimes m} \in K_T(M_{\mathbb{D}})[[t]] \otimes \mathbb{Q}$$

$$\Rightarrow \text{index}_k(e^{tE^{\mathbb{D}}V})(u) = J(u) \sum_{\eta \in \Pi} \Phi_t(u)^{\eta}, \quad \Phi_t(u) = u e^{-t \nabla \chi_V(u)}$$

(formal diffeomorphism of T)

$$\stackrel{\text{Poisson}}{=} \frac{J(u)}{|T_{\ell}|} \sum_{\mathcal{J} \in T_{\ell}} \frac{\delta_{\mathcal{J}_t}(u)}{\det(1 - t \nabla^2 \chi_V(u))}$$

where $\Phi_t(\mathcal{J}_t) = \mathcal{J}$ and $T_{\ell} = \ell^{-1}(1)/\Pi \subseteq T$.

Remarks on $K_T(M_D)$

$$\iota^*: K_T(M_D) \longrightarrow K_T(M_D^{T \times S^1_{\text{rot}}}) = K_T(\Pi) = \text{Fun}(\Pi, R(T))$$

is injective, image consists of

$$\psi: \Pi \rightarrow R(T) \quad \text{st} \quad (1 - e^\alpha)^m \text{ divides } \Delta_{\alpha^\vee}^m \psi$$

for each $m > 0$ and root α of G .

finite diff. operator

[Lam - Schilling - Shimozono], [Kostant - Kumar]

$$(K_T(M_D) \cong \prod_{\gamma \in \Pi} R(T) \text{ as } R(T)\text{-mod.}, \text{ Schubert basis})$$

Non-abelian localization

Recall index_k takes $\psi = [(\mathcal{E}, \mathcal{Q})] \in K_T(M)_{gc}$ to the L^2 -index of

$$D = \not{D}_{X, \xi}^{\mathcal{L}} \hat{\otimes} 1 + 1 \hat{\otimes} \mathcal{Q}_X$$

Let $D_s = D + s c(k)$, $k = \omega^{-1}(d \|\mu\|^a) / (1 + \|\mu\|^a)^{1/2}$

By considering $s \rightarrow \infty$ get

Thm [L] $\text{index}_k(\psi) = \sum_{v \in \text{Crit}(\|\mu\|^a)} \text{index}_{\ell, v}(\psi|_{C_v})$

Rmk $C_v \cap \text{supp}(\xi)$ compact, sum over v infinite.

c.f.

[Witten], [Bismut-Lebeau],

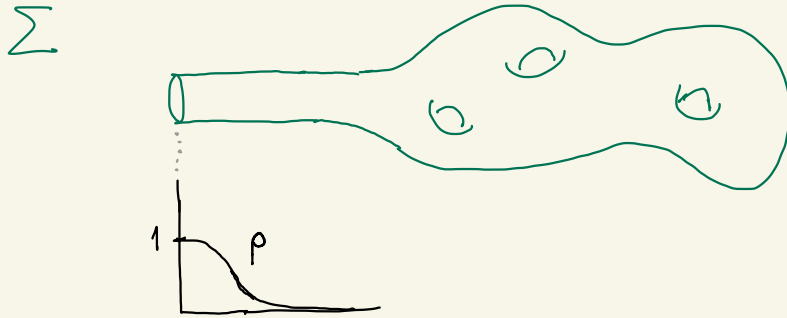
[Tian-Zhang], [Ma-Zhang],

[Paradan], [Hochs-Song],

[L-Song], ...

Remarks on κ

For M_Σ



the flow of $\rho \partial_\theta \in \mathcal{X}(\Sigma)$
induces a flow on M_Σ
gives the flow of κ
(up to reparametrization)

For $\Sigma = \mathbb{D}$, orbits close, S'_{rot} -action. The sum over $\Pi = \mathcal{M}^{\text{T} \times S'_{\text{rot}}}$
was the non-abelian localization formula in that case.

Main theorem

Under further assumptions on $\psi \in K_T(M)_{gc}$ one gets both a Kirillov-Berline-Vergne formula & Atiyah-Bott-Segal-Singer formula.

We'll just state the latter for the special case $M = M_\Sigma$, $\psi = E^\Sigma V$.

Thm [L] $\text{index}_\ell(E^\Sigma V) = \frac{1}{J} \sum_{\mathcal{J} \in T_\ell^{\text{reg}}} \left(\frac{J^a}{|T_\ell| \det(1 - t \nabla^2 \chi_V)} \right)^{1 - \text{genus}(\Sigma)} \int_{\mathcal{Y}_t}$

where $\Phi_t(\mathcal{Y}_t) = \mathcal{Y}$, $\Phi_t(u) = u e^{-t \nabla \chi_V(u)}$.

Rmk This is a gauge theoretic version of Teleman-Woodward's algebro-geometric result.

Rmk Similar for any number of boundary components & other Atiyah-Bott classes.

Proof outline

$$\text{index}_\ell(E^\pm V) = \sum_{v \in \text{Crit}(\|\mu\|^2)} \text{index}(D_v + f_v c(\beta_v))$$

↙ Tian-Zhang / Braverman - type deformation
of $D_v = D|_{U_v}$, with
 $f_v \rightarrow \infty$ at ∂U_v

Work of Braverman & Paradan - Vergne \rightarrow index formula for $\xi \in \mathfrak{t}$ small:

$$\text{index}(D_v + f_v c(\beta_v))(\exp \xi) = \int_{U_v} \hat{A}(U_v, \xi) \text{Ch}(E^\pm V \otimes \mathcal{L}, \xi) P_v(\xi) \text{Th}(\xi)$$

Non-abelian localization in ordinary cohomology (Paradan)

↖ localizing form
(generalized coeffs)

$$\text{index}_\ell(E^\pm V)(\exp \xi) = \int_X \hat{A}(X, \xi) \text{Ch}(E^\pm V \otimes \mathcal{L}, \xi) \text{Th}(\xi)$$

Compute $\text{Ch}(E^\pm V, \xi)$ (families index thm) & $\eta^* \text{Ch}(E^\pm V, \xi) = \text{Ch}(E^\pm V, \xi) + \nabla_\eta \chi_V(\exp \xi)$

$$\Rightarrow \eta^* \text{Ch}(e^{tE^\pm V} \otimes \mathcal{L}, \xi) = e^{\ell(\eta, \xi) + t \nabla_\eta \chi_V(\exp \xi)} \text{Ch}(e^{tE^\pm V} \otimes \mathcal{L}, \xi)$$

$$\int_X \square = \sum_{\eta \in \Pi} \int_{X_0} \eta^* \square \quad \text{Poisson summation. Evaluate resulting equiv. integral over } X/\Pi.$$

Thanks !