

Delocalized invariants for proper actions of Lie groups

Based on joint work with P. Piazza, Y. Song and X. Tang

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Motivation: Numerical invariants from index classes

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$$H^\bullet(\Gamma) \rightarrow HC^\bullet(\mathbb{C}[\Gamma]), \quad \varphi \mapsto \tau_\varphi$$

D invariant Dirac operator $\implies \text{Ind}(D) \in K_0(C_r^*(\Gamma))$

If τ_φ extends to a subalgebra $\mathbb{C}[\Gamma] \subset \mathcal{A} \subset C_r^*(\Gamma)$ we can pair:

$$\langle \text{Ind}(D), \tau_\varphi \rangle = \int_{M/\Gamma} AS(D) \wedge \nu^* \varphi.$$

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- ▶ **Secondary invariants:** higher APS index theorem (Leichtnam–Piazza, . . .). Assume D_∂ is L^2 -invertible:

$$\langle \text{Ind}(D), \tau_\varphi \rangle = \int_{M/\Gamma} AS(D) \wedge \nu^* \varphi - \frac{1}{2} \eta_\varphi(D_\partial).$$

$\eta_\varphi(D_\partial)$ higher η -invariant.

$$\eta_0(D) := \frac{1}{\sqrt{\pi}} \int_0^\infty \text{Tr} \left(D e^{-tD^2} \right) \frac{dt}{\sqrt{t}}$$

\implies higher ρ -invariants.

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- ▶ Instead of discrete groups consider Lie groups (connected, real reductive)
- ▶ Allow for general proper actions, instead of free actions
- ▶ Consider *delocalized cyclic cocycles*.
- ▶ N.B. $HC^\bullet(C_c^\infty(G))$ is a module over

$$C_{\text{inv}}^\infty(G) := \{f \in C^\infty(G), f(hgh^{-1}) = f(g)\}.$$

A cocycle is called *delocalized* if its image in $HC^\bullet(C_c^\infty(G))_{m_e}$ is zero, where $m_e := \{f, f(e) = 0\}$.

- ▶ The Plancherel trace is *localized at the identity*:

$$\tau_e(f) := f(e).$$

Paolo's questions

- ▶ can we define a delocalized eta invariant $\eta_g(D)$ using the heat kernel?
- ▶ if C is a smoothing perturbation can we define $\eta_g(D + C)$?
- ▶ in particular, can we define the delocalized eta invariant of a PSC metric g and of a G -equivariant homot. equivalence f ?
- ▶ is there a delocalized APS index theorem ?
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The answer to all these questions is: **Yes!**

Outline

Delocalized η -invariants

APS index theorem

Perturbations

Higher versions

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- ▶ $g \mapsto \|g\|$ riemannian distance from eK to gK in G/K .
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- ▶ $\Xi(g)$ is Harish-Chandra's Ξ -function $\Xi(g) := \int_K a(kg)^p dk$
- ▶ Let $g \in G$ semisimple. Orbital integral

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- ▶ N.B. Other versions: Harish-Chandra algebra $\mathcal{C}(G)$, rapid decay $H_L^\infty(G)$.

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 - ▶ **metric**: K -invariant metrics on S and \mathfrak{p} via $TX \cong G \times_K (TS \oplus \mathfrak{p})$.
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 - ▶ **Spin^c-structure**: K -invariant Spin^c-structures on S and \mathfrak{p}
- ▶ **Dirac operator**

$$D = D_{G,K} \hat{\otimes} 1 + 1 \hat{\otimes} D_S,$$

$D_{G,K}$ Spin^c-Dirac operator on $L^2(G) \otimes S_{\mathfrak{p}}$.

Delocalized traces

- ▶ Subalgebras of the Roe algebra $C^*(X)^G$:

$$\begin{aligned}\mathcal{A}_G^c(X) &:= (C_c^\infty(G) \hat{\otimes} \Psi^{-\infty}(S))^{K \times K} \\ &\cong \{\Phi : G \rightarrow \Psi^{-\infty}(S), \text{ compact support, } K \times K \text{ invariant.}\end{aligned}$$

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- ▶ delocalized trace $\tau_g^X : \mathcal{A}_G^\infty(X) \rightarrow \mathbb{C}$ given by:

$$\tau_g^X(\Phi) := \tau_g(\text{Tr}_S(\Phi)).$$

Delocalized η -invariants

X is of bounded geometry and $\mathcal{A}_G^\infty(X) \subset C^*(X)^G$ is closed under holomorphic functional calculus:

$$e^{-tD^2} = \frac{1}{2\pi i} \int_C \frac{e^{-t\lambda}}{\lambda - D^2} d\lambda \in \mathcal{A}_G^\infty(X).$$

Follows from analysis of the resolvent

$$\lambda - D^2 = B_\lambda + C_\lambda, \quad B_\lambda \in \Psi_c^{-2}(X), \quad C_\lambda \in \mathcal{A}_G^\infty(M),$$

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Theorem (Piazza–P.–Song–Tang)

The following integral converges:

$$\eta_g(D) := \frac{1}{\sqrt{\pi}} \int_0^\infty \tau_g^X(D e^{-tD^2}) \frac{dt}{\sqrt{t}}$$

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N.B: No invertibility or gap assumptions!

About the proof

Large time behaviour: Recall $D^2 = D_{G,K}^2 + D_S^2$ and decompose

$$L^2(X) = \bigoplus_{\lambda_i \in \sigma(D_S)} \left[L^2(G) \otimes \mathfrak{p} \otimes L^2(S)_{\lambda_i} \right]^K.$$

Moscovici–Stanton give estimates

$|\tau_g^{G/K}(D_{G,K} e^{-D_{G,K}^2})| \leq C_1 e^{-C_2/t} t^{-3/2}$ resulting in convergence of

$$\frac{1}{\sqrt{\pi}} \int_1^\infty \tau_g^X(D e^{-tD^2}) \frac{dt}{\sqrt{t}}.$$

About the proof

Short time behaviour: rewrite (Hochs–Wang)

$$\tau_g^X(De^{-tD^2}) = \int_X \int_G c_G(g)c_X(gx)\text{tr}(k_t(x, gx))dgdx,$$

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Let $W \subset X$ open subset containing X^g .

- ▶ Gaussian estimates for the heat kernel gives exponential decay of

$$\int_{X \setminus W} \int_G c_G(g)c_X(gx)\mathrm{tr}(k_t(x, gx))dgdx, \quad \text{as } t \downarrow 0.$$

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- ▶ An argument of Zhang shows that

$$\frac{1}{\sqrt{t}} \int_W \int_G c_G(g)c_X(gx)\mathrm{tr}(k_t(x, gx))dgdx = O(1) \quad t \downarrow 0.$$

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Using b -calculus with ϵ -bounds on the slice $S \subset Y$:

$$0 \longrightarrow \mathcal{A}_G^c(Y) \longrightarrow {}^b\mathcal{A}_G^c(Y) \xrightarrow{I} {}^b\mathcal{A}_{G,\mathbb{R}}^c(\text{cyl}(\partial Y)) \longrightarrow 0.$$

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Using Lafforgue's algebra:

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$$(D^+ Q^b = 1 - {}^b S_-, \quad Q^b D^{-1} = 1 - {}^b S_+)$$

$$P_Q^b := \begin{pmatrix} {}^b S_+^2 & {}^b S_+(I + {}^b S_+)Q^b \\ {}^b S_- D^+ & I - {}^b S_-^2 \end{pmatrix}$$

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gives an index class

$$\text{Ind}_\infty(D) := [P_Q^b] - [e_1] \in K_0(\mathcal{A}_G^\infty(Y)) = K_0(C^*(Y_0 \subset Y)^G)$$

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$$V(D) := \begin{pmatrix} e^{-D^- D^+} & e^{-\frac{1}{2} D^- D^+} \left(\frac{I - e^{-D^- D^+}}{D^- D^+} \right) D^- \\ e^{-\frac{1}{2} D^+ D^-} D^+ & I - e^{-D^+ D^-} \end{pmatrix} \in M_2({}^b\mathcal{A}_G^\infty(Y))$$

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Relative index class

$\text{Ind}_\infty(D, D_\partial) \in K_0({}^b\mathcal{A}_G^\infty(Y_0), {}^b\mathcal{A}_{G, \mathbb{R}}^\infty(\text{cyl}(\partial Y)))$:

$$(V(D), e_1, q_t), \quad t \in [1, +\infty], \quad \text{with } q_t := \begin{cases} V(tD_{\text{cyl}}) & \text{if } t \in [1, +\infty) \\ e_1 & \text{if } t = \infty \end{cases}$$

Relative cyclic cocycles

- ▶ Recall the delocalized trace:

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- ▶ Using the b -trace $\mathrm{Tr}_S^b(k) := \int_S^b k(x, x) dx$,

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- ▶ 1-cocycle on ${}^b\mathcal{A}_{G,\mathbb{R}}^\infty(\mathrm{cyl}(\partial Y))$:

$$\sigma_g^{\partial Y}(A_0, A_1) := \frac{i}{2\pi} \int_{\mathbb{R}} \tau_g^{\partial Y}(\partial_\lambda I(A_0, \lambda) \circ I(A_1, \lambda)) d\lambda,$$

where the indicial family of $A \in {}^b\mathcal{A}_{G,\mathbb{R}}^\infty(\mathrm{cyl}(\partial Y))$, denoted $I(A, \lambda)$, appears.

Index pairings

- ▶ Recall Melrose formula for the b -trace:

$${}^b\mathrm{Tr}_S([A_0, A_1]) = \frac{i}{2\pi} \int_{\mathbb{R}} \mathrm{Tr}_{\partial S}(\partial_\lambda(A_0, \lambda) \circ I(A_1, \lambda)) d\lambda$$

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- ▶ This implies

$$\begin{pmatrix} (b+B) & -I^* \\ 0 & -(b+B) \end{pmatrix} \begin{pmatrix} \tau_g^{Y,r} \\ \sigma_g^{\partial Y} \end{pmatrix} = 0,$$

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$$\begin{pmatrix} (b+B) & -I^* \\ 0 & -(b+B) \end{pmatrix} \begin{pmatrix} \tau_g^{Y,r} \\ \sigma_g^{\partial Y} \end{pmatrix} = 0,$$

- ▶ and leads to

$$\langle \tau_g^Y, \mathrm{Ind}_\infty(D) \rangle = \langle (\tau_g^{Y,r}, \sigma_g), \mathrm{Ind}_\infty(D, D_\partial) \rangle.$$

The delocalized APS Theorem

Theorem (Piazza–P.–Song–Tang)

Assume that D_∂ is L^2 -invertible.

$$\langle \tau_g^Y, \text{Ind}_\infty(D) \rangle = \int_{Y_0^g} c^g \text{AS}_g(D_0) - \frac{1}{2} \eta_g(D_\partial Y),$$

where

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- ▶ Closed manifold case is due to Hochs–Wang.

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$$\langle (\tau_g^{Y,r}, \sigma_g), \text{Ind}_\infty(D, D_\partial) \rangle = \tau_{g,s}^{Y,r}(e^{D^2}) + \int_1^\infty \sigma_g^{\partial Y}([\dot{q}_t, q_t], q_t) dt$$

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- ▶ Pairing on the boundary with $\sigma_g^{\partial Y}$ produces a long complicated expression that we show that this equals $\eta_g(D)$.

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For the signature operator D_∂ is not invertible.

What if D_∂ is not L^2 -invertible?

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- ▶ general perturbation $D + C$ with $C \in \mathcal{A}_G^c(X)$. (always exists!)

Global perturbation on the slice

- ▶ Recall $D = D_{G,K} \otimes 1 + \gamma \otimes D_S$. The Connes–Kasparov correspondence gives

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- ▶ APS-type index theorem, by letting $\vartheta \downarrow 0$:

Theorem (Piazza–P.–Song–Tang)

Assume G is of equal rank, i.e., has discrete series. Then

$$\langle \text{Ind}_\infty(D_\vartheta), \tau_g^Y \rangle = \int_{Y_0^g} c^g \text{AS}_g(D_0) - \frac{1}{2} (\eta_g(D_{\partial Y}) + \langle \text{Ind}_D(\ker(D_S)), \tau_g \rangle)$$

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General perturbations

- ▶ Previous cases work under a gap assumption. We now consider

$$D + C, \quad C \in \mathcal{A}_G^c(X), \quad D + C \text{ is } L^2\text{-invertible.}$$

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Theorem (Piazza–P.–Song–Tang)

Assume that $D_\partial + B_\partial^\infty$ is L^2 -invertible.

$$\langle \text{Ind}_\infty(D, D_\partial + B_\partial^\infty), (\tau_g^{Y,r}, \sigma_g^{\partial Y}) \rangle = \int_{Y_0^g} c^g \text{AS}_g(D_0) - \frac{1}{2} \eta_g(D_\partial + B_\partial^\infty)$$

Higher versions

- ▶ Song–Tang: For any cuspidal parabolic subgroup $P = MAN \subset G$ and $g \in G$ semisimple there exists a cyclic cocycle Φ_g^P of degree $\dim(A)$ on $\mathcal{L}_t(G)$.

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- ▶ \Rightarrow cyclic cocycle $\Phi_{Y,g}^P$ on $\mathcal{A}_G^\infty(Y)$.
- ▶ Higher APS index theorem:

Theorem (Piazza–P.–Song–Tang)

Assume that D_∂ is L^2 -invertible.

$$\langle \Phi_{Y,g}^P, \text{Ind}_\infty(D) \rangle = \int_{(Y_0/AN)^g} c_{Y_0/AN}^g \text{AS}(Y_0/AN)_g - \frac{1}{2} \eta_g(D_{\partial Y_M}),$$

where $Y_M := M \times_{K \cap M} S$ is the slice decomposition of Y/AM with its M -action.

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- ▶ If $g \in G$ is not elliptic, any element of its conjugacy class C_g will act fixed point free on any proper G -manifold W :

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- ▶ $\mathbf{f} : X_1 \rightarrow X_2$ G -homotopy invariant. Fukumoto: there exists a perturbation $B_{\mathbf{f}}$ making the signature operator on $X := X_1 \sqcup (-X_2)$:

$$\rho_g(\mathbf{f}) := \eta_g(D_X^{\text{sign}} + B_{\mathbf{f}}).$$

Bordism invariant by the APS index theorem.

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